

# Theoretical Elementary Particle Physics

## Exercise 2

May 22, 2015

-to be handed in on Friday, 5th of June

### 1. Feynman parametrization (25 points)

**(a)(10 points)** Using standard integration techniques prove the identity :

$$\frac{1}{A B C} = 2 \int_0^1 dx \int_0^{1-x} dz \frac{1}{[A + (B - A)x + (C - A)z]^3} \quad (1)$$

**(b)(15 points)** Prove the following formula

$$\frac{1}{A_1 A_2 A_3 \dots A_n} = (n-1)! \int_0^1 \dots \int_0^1 dx_1 dx_2 \dots dx_n \frac{\delta(x_1 + x_2 + \dots + x_n - 1)}{(x_1 A_1 + x_2 A_2 + x_3 A_3 + \dots + x_n A_n)^n} \quad (2)$$

( **Hint:** Use mathematical induction. Useful relation is

$$\frac{1}{A_1 A_2^n} = \int_0^1 dx_1 \int_0^1 dx_2 \delta(x_1 + x_2 - 1) \frac{n x_2^{n-1}}{(x_1 A_1 + x_2 A_2)^{n+1}} \quad (3)$$

If you use it prove it by performing several derivatives on  $(A_1 A_2)^{-1}$  with respect to  $A_2$  and looking for a pattern or, alternatively, by means of the mathematical induction.)

### 2. 1-loop correction to the propagator in $\phi^4$ theory (35 points)

**(a)(30 points)** Using Feynman rules for  $\phi^4$  theory (see exercise!) write the amplitude corresponding to the Feynman diagram in figure 1. The symmetry factor of this diagram is 2. Use dimensional regularization. The angular part of the d-dimensional integral equals  $\frac{2\pi^{d/2}}{\Gamma(\frac{d}{2})}$ . In the spatial part it is useful to make a substitution of the integration variable by introducing  $x = \frac{\Delta}{l_E^2 + \Delta}$ . In this simplified example, which does not require usage of Feynman parametrization,  $\Delta \equiv m^2$ . Another useful identity is :

$$\int_0^1 dx x^{\alpha-1} (1-x)^{\beta-1} = \frac{\Gamma(\alpha)\Gamma(\beta)}{\Gamma(\alpha+\beta)} \quad (4)$$

Take  $d = 4 - \epsilon$  and expand the amplitude in powers of  $\epsilon$ . Consider only  $\epsilon^{-1}$  (divergent part) and  $\epsilon^0$  parts since other terms will vanish as you go to  $\epsilon \rightarrow 0$  ( $d \rightarrow 4$ ) limit. For the expansion of the  $\Gamma$  function near  $z = -1$  pole use :

$$\Gamma(-1 + \frac{\epsilon}{2}) \approx -\frac{2}{\epsilon} + \gamma - 1 \quad (5)$$

where  $\gamma$  stands for the Euler-Mascheroni constant introduced in the exercises.

**(b)(5 points)** After extracting the  $\epsilon^{-1}$  part in dimensional regularization, absorb this divergence into the renormalization constants  $\delta_Z$  and  $\delta_m$  by adding the amplitude of the diagram and the counterterm in figure 2 and requiring cancellation of infinite parts. Use the  $\overline{\text{MS}}$  renormalization scheme, i.e. absorb only the divergent part in  $\delta_m$  and  $\delta_Z$ .

**Solution :**  $\delta_Z = 0$ ,  $\delta_m = \frac{\lambda m^2}{16\pi^2\epsilon}$

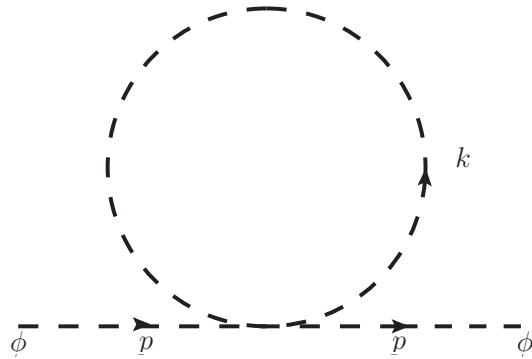


Figure 1: One loop correction to the propagator in  $\phi^4$  theory with 4-momenta labelled.



Figure 2: Counterterm that absorbs infinite contributions for diagram in figure 1. The Feynman rule for this counterterm is  $i(p^2\delta_Z - \delta_m)$ .

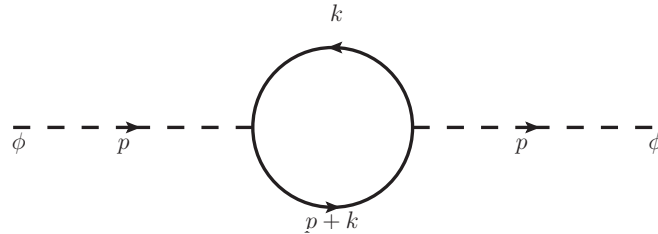


Figure 3: One loop correction to the propagator in Yukawa theory (interaction Lagrangian is  $-g\bar{\psi}\psi\phi$ ) with 4-momenta labelled.

### 3. 1-loop correction to the propagator in Yukawa theory (40 points)

One can study the interaction between a fermion field  $\psi$  and scalar field  $\phi$  which

is described by the Lagrangian term  $-g\bar{\psi}\psi\phi$ ,  $g$  being the coupling constant. The 1-loop correction to the scalar propagator is presented in figure 3. The amplitude may be expressed in the following form

$$-g^2 \int \frac{d^d k}{(2\pi)^d} \text{Tr} [(k + \not{p} + m)(k + m)] \frac{1}{((k + p)^2 - m^2)(k^2 - m^2)} \quad (6)$$

where  $m$  now represents the mass of the fermion propagating in the loop.

**(a)(30 points)** Starting from equation (6), perform dimensional regularization using the procedure described in the previous problem. Note that in this case Feynman parametrization should be employed. When shifting the integration 4-momentum variable  $k$  to  $l = k + xp$  note that the terms in the numerator which are linear in  $l$  should be dropped because the integral of an asymmetrical function on a symmetrical interval vanishes. Show that the amplitude may be written in the form :

$$(d-1) \frac{4ig^2}{(4\pi)^{d/2}} \int_0^1 dx \frac{\Gamma(1 - \frac{d}{2})}{\Delta^{1 - \frac{d}{2}}} \quad (7)$$

where  $\Delta = x(x-1)p^2 + m^2$  and  $x$  is a Feynman parameter.

**(b)(10 points)** Take  $d = 4 - \epsilon$  and expand the factors containing  $\epsilon$ . Use again equation (5) to expand  $\Gamma$  function near pole  $z = -1$ . After identifying the divergent part, using the counterterm from figure 2, obtain the expressions for renormalization constants  $\delta_m$  and  $\delta_Z$ . Use the MS renormalization scheme, i.e. absorb only the divergent part in  $\delta_m$  and  $\delta_Z$ . **Solution :**  $\delta_Z = -4 \frac{g^2}{(4\pi)^2 \epsilon}$ ,  $\delta_m = -\frac{24g^2 m^2}{(4\pi)^2 \epsilon}$